

## Symmetry reduction in the electronic structure of heavily overdoped Pb-Bi2201 detected by ARPES

Y. Miyai<sup>a</sup>, T. Kurosawa<sup>b</sup>, M. Oda<sup>c</sup>, M. Arita<sup>d</sup>, S. Ideta<sup>d</sup>, and K. Shimada<sup>d</sup>

<sup>a</sup>*Graduate School of Science, Hiroshima University, Kagamiyama 1-3-1, Higashi-Hiroshima 739-8526, Japan*

<sup>b</sup>*Fuculty of Science and Engineering, Muroran Institute of Technology, 27-1 Mizumoto-cho, Muroran, Hokkaido, 050-8585, Japan*

<sup>c</sup>*Department of Physics, Hokkaido University, Sapporo 060-0809, Japan*

<sup>d</sup>*Hiroshima Synchrotron Radiation Center, Hiroshima University, Kagamiyama 1-3-1, Higashi-Hiroshima 739-8526, Japan*

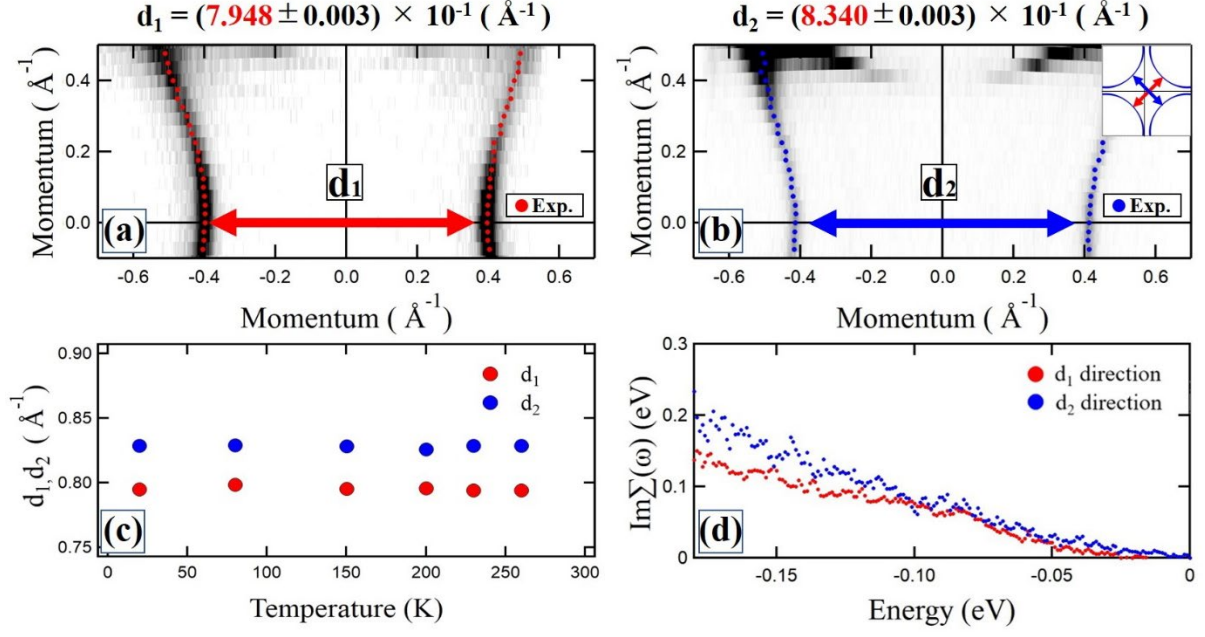
**Keywords:** High- $T_C$  superconductor, ARPES

High transition-temperature ( $T_C$ ) cuprate superconductors have attracted much interest since their discovery in 1986 for their high superconducting transition temperature as well as unusual physical properties such as a pseudogap state and a nematic phase [1,2]. Superconductivity occurs in the  $\text{CuO}_2$  plane and the  $\text{Cu } 3d_{x^2-y^2}$  state forms the Fermi surface. One can expect a four-fold symmetry of the Fermi surface because the  $\text{CuO}_2$  plane is tetragonal. However, symmetry reduction of the electronic structure has been reported in Bi-based cuprate superconductors recently [2,3]. Similar symmetry reduction or *nematicity* was found in the Fe-based superconductor, the electronic states break the four-fold rotational symmetry of the lattice [4]. Although nematicity in electron liquids has attracted much interest, the rotational symmetry breaking in the electronic states for high- $T_C$  cuprate superconductors has not been clarified yet.

Here, we have examined the symmetry of the electronic structure of heavily overdoped  $(\text{Bi,Pb})_2\text{Sr}_2\text{CuO}_{6+\delta}$  (Pb-Bi2201) ( $T_C = 6 \text{ K}$ ) using high-resolution angle-resolved photoemission spectroscopy (ARPES). We selected Bi2201 because it has a single  $\text{CuO}_2$  plane and one can exclude the complexity of the electronic structure such as bilayer splitting derived from plural  $\text{CuO}_2$  planes. Furthermore, the heavily overdoped sample is favorable because the pseudogap is vanishing and one can clearly detect the Fermi surface in the normal state.

Figures 1(a) and (b) show the Fermi surface measured at  $h\nu = 22 \text{ eV}$ ,  $s$ -polarization, and  $T = 20 \text{ K}$  along two orthogonal nodal directions. We found a clear difference in the nodal distances;  $d_1 = (7.948 \pm 0.003) \times 10^{-1} \text{ \AA}^{-1}$  (Fig. 1(a)) and  $d_2 = (8.340 \pm 0.003) \times 10^{-1} \text{ \AA}^{-1}$  (Fig. 1(b)). These distances are temperature independent for the temperature range from  $T = 20 \text{ K}$  up to  $260 \text{ K}$  (Fig. 1(c)). In addition, we have conducted MDC analysis along these two nodal directions and find different quasiparticle lifetime broadening as shown in Fig. 1(d), which is persistent up to  $260 \text{ K}$ . The situation is similar for the samples with  $T_C = 4, 6$ , and  $10 \text{ K}$ . Our results indicate the reduction of the four-fold symmetry in the electronic states.

In the heavily overdoped region, a charge density wave (CDW) is observed by resonant inelastic X-ray scattering and its onset is well above  $250 \text{ K}$  [5]. In addition, previous Raman scattering measurements suggest the Pomeranchuk instability and the nematic phase [6]. The previous reports are in line with our findings. Further study is required to reveal how symmetry reduction in the electronic state occurs in the heavily overdoped region. To investigate this issue, we are planning to examine the electronic structure at the endpoint of the hole concentration in the pseudogap state and at the onset of the Lifshitz transition that occurs between the hole-like and electron-like Fermi surfaces.



**FIGURE 1.** (a), (b) Fermi surface observed at  $h\nu = 22$  eV,  $s$ -polarization and  $T = 20$  K along the nodal directions shown in the inset of panel (b). The inset displays the Fermi surface centered at the  $\bar{\Gamma}$  point. Red and blue dots indicate the Fermi surface points determined from the peaks of the momentum distribution curves.  $d_1$  and  $d_2$  indicate the distance between nodes. (c) Temperature dependence of  $d_1$ ,  $d_2$ . (d)  $\text{Im}\Sigma$  of the self-energy along  $d_1$  and  $d_2$  nodal directions.

1. M. Hashimoto *et al.*, *Nature Physics* **10**, 483-495 (2014).
2. S. Nakata *et al.*, *npj Quantum Materials* **6**, 86 (2021).
3. Y. He *et al.*, *Scientific Reports* **7**, 8059 (2017).
4. H. C. Xu *et al.*, *Physical Review Letters* **117**, 157003 (2016).
5. Y. Y. Peng *et al.*, *Nature Materials* **17**, 697 (2018).
6. N. Auvray *et al.*, *Nature Communications* **10**, 5209 (2019).