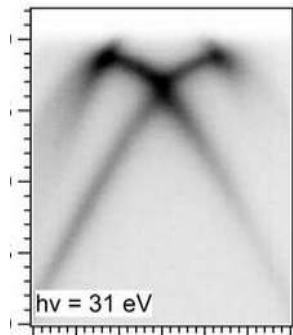


# Unscrambling orbital character and spin: The importance of the final state

CORPES17

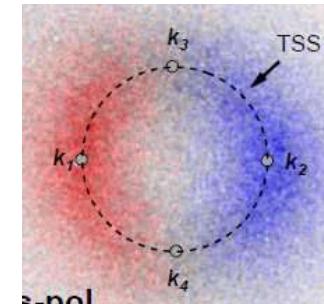
Hiroshima



6. July 2017

Friedel Reinert

[reinert@physik.uni-wuerzburg.de](mailto:reinert@physik.uni-wuerzburg.de)



# Acknowledgments

Würzburg Group: SARPES experiments

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- **E. Krasovskii (San Sebastian)**

- **J. Braun, H. Ebert, J. Minar (LMU, UWB)**

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- **T. Balasubramanian, M. Leandersson (MAXLAB)**

MPI Halle, Forschungszentrum Jülich: Momentum Microscope

- **Christian Tusche**

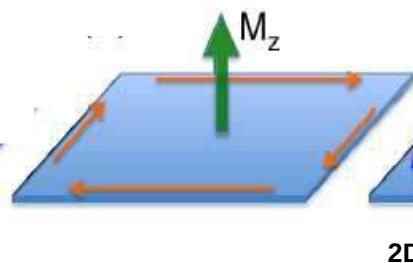


Bundesministerium  
für Bildung  
und Forschung

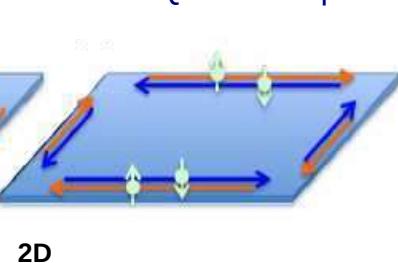


# Motivation: Spin-split electronic structure

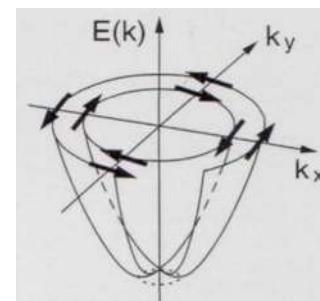
Quantum Anomalous Hall effect



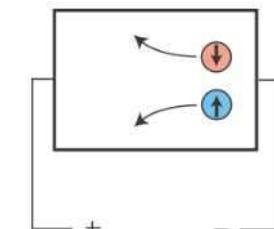
Quantum Spin Hall effect



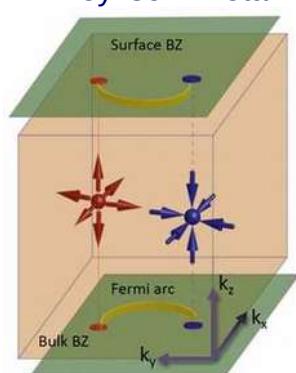
Rashba effect



Spin-Hall effect



Weyl semimetal



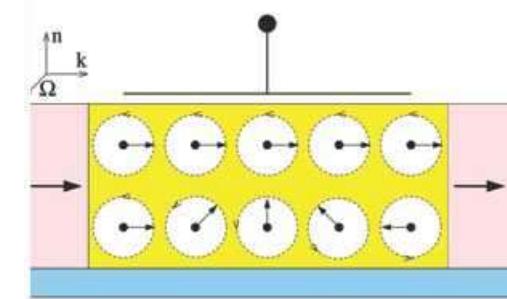
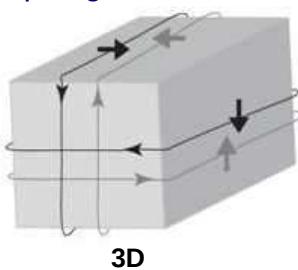
## Spin-Orbit Coupling (SOC)



spin-polarization of electronic states  
→ origin of non-trivial topology

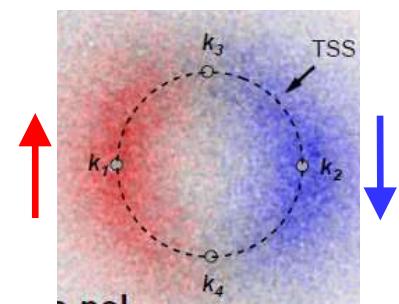


Topological insulator  
Topological Kondo insulator



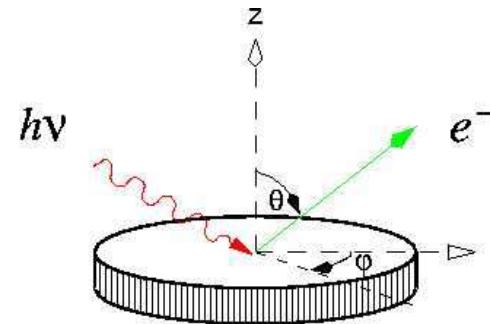
Spin transistor

direct spectroscopy by spin-  
and angle-resolved  
photoemission

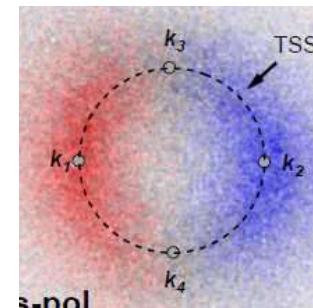


# Outline

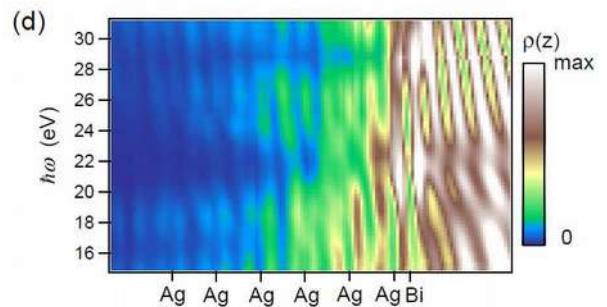
ARPES on systems with strong spin-orbit coupling



Complications by initial state  
→ linear dichroism

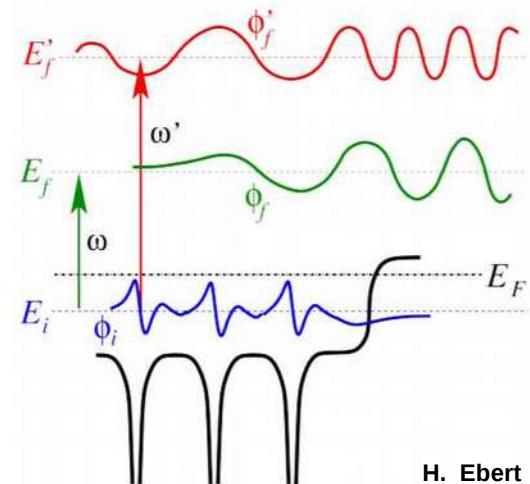
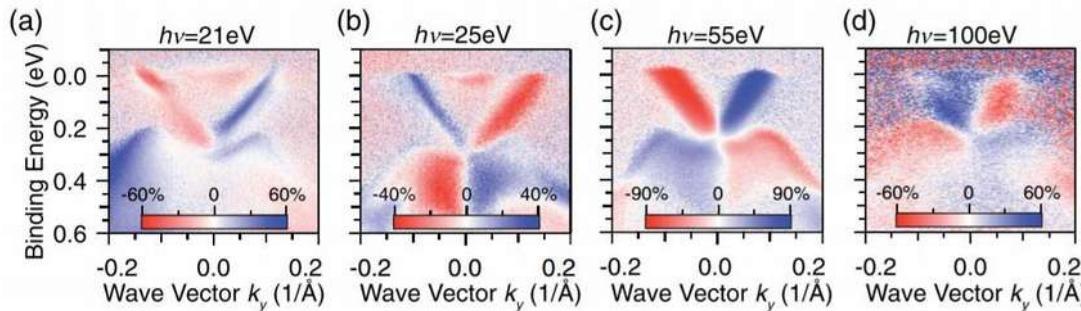


Importance of the final state  
→ photoelectron polarization



# How to address intrinsic spin properties by photoemission?

## Circular dichroism (indirectly)

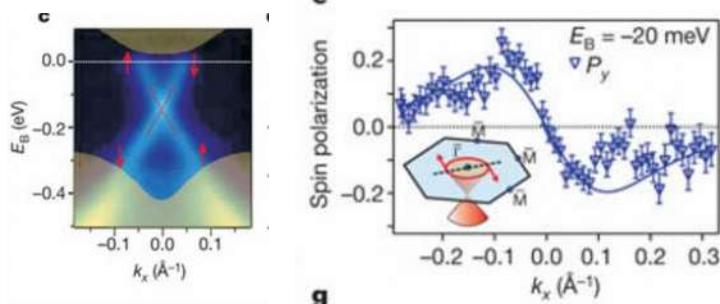


H. Ebert

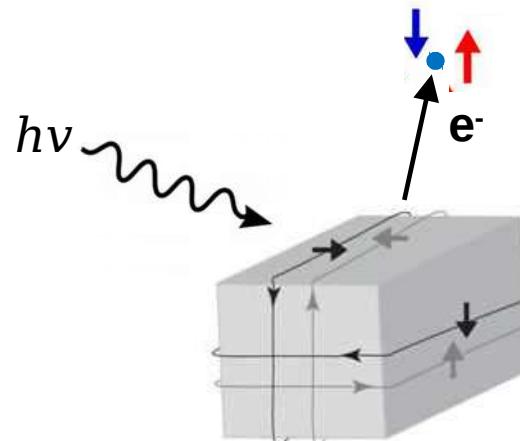
Scholz et al., Phys. Rev. Lett. **110**, 216801 (2013)

→ „final state effect“, dichroism depends on excitation energy

## Photoelectron spin polarization



Hsieh et al., Nature **460**, 1101-1105 (2009)



→ also dependence on photoemission process?

**VIEWPOINT**

# Can spin-polarized photoemission measure spin properties in condensed matter?

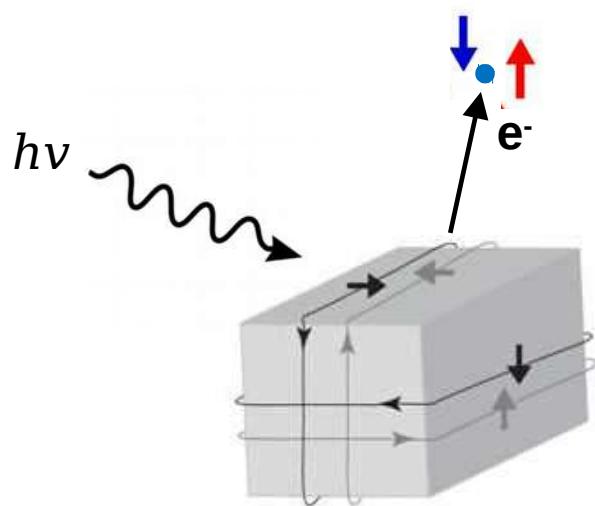
**Jürg Osterwalder**

Physik-Institut, Universität Zürich, Winterthurerstraße 190, 8057 Zürich, Switzerland

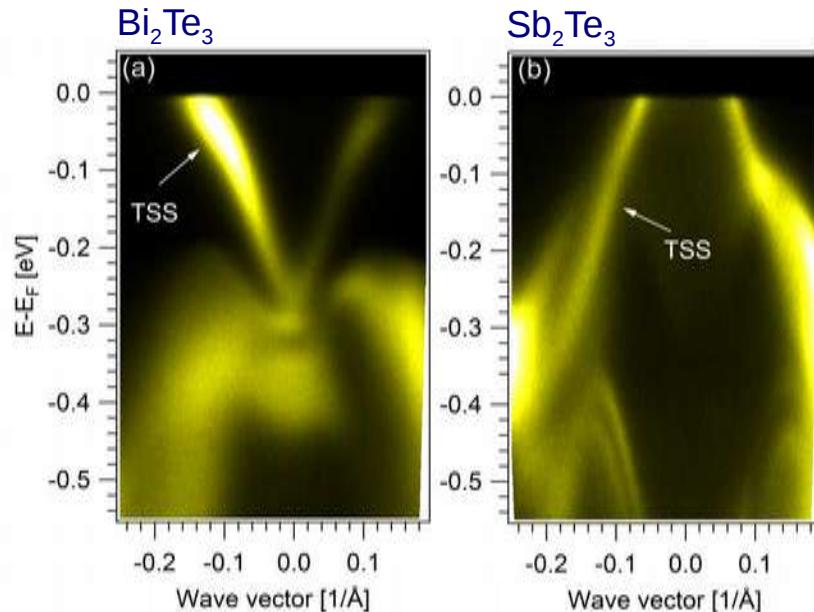
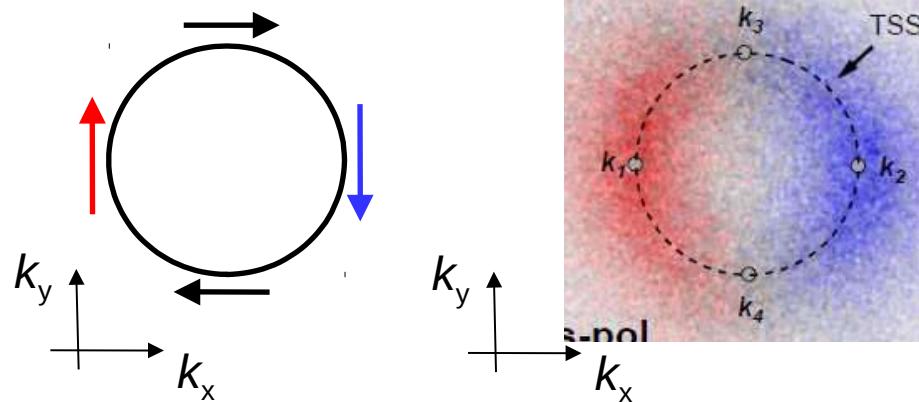
The question posed in the title of this viewpoint article cannot be answered conclusively at this point.

these ‘intrinsic’ spin polarizations can be expected to be stable with respect to variations in the photon energy and photon polarization direction

# Spin- and angle-resolved photoelectron spectroscopy



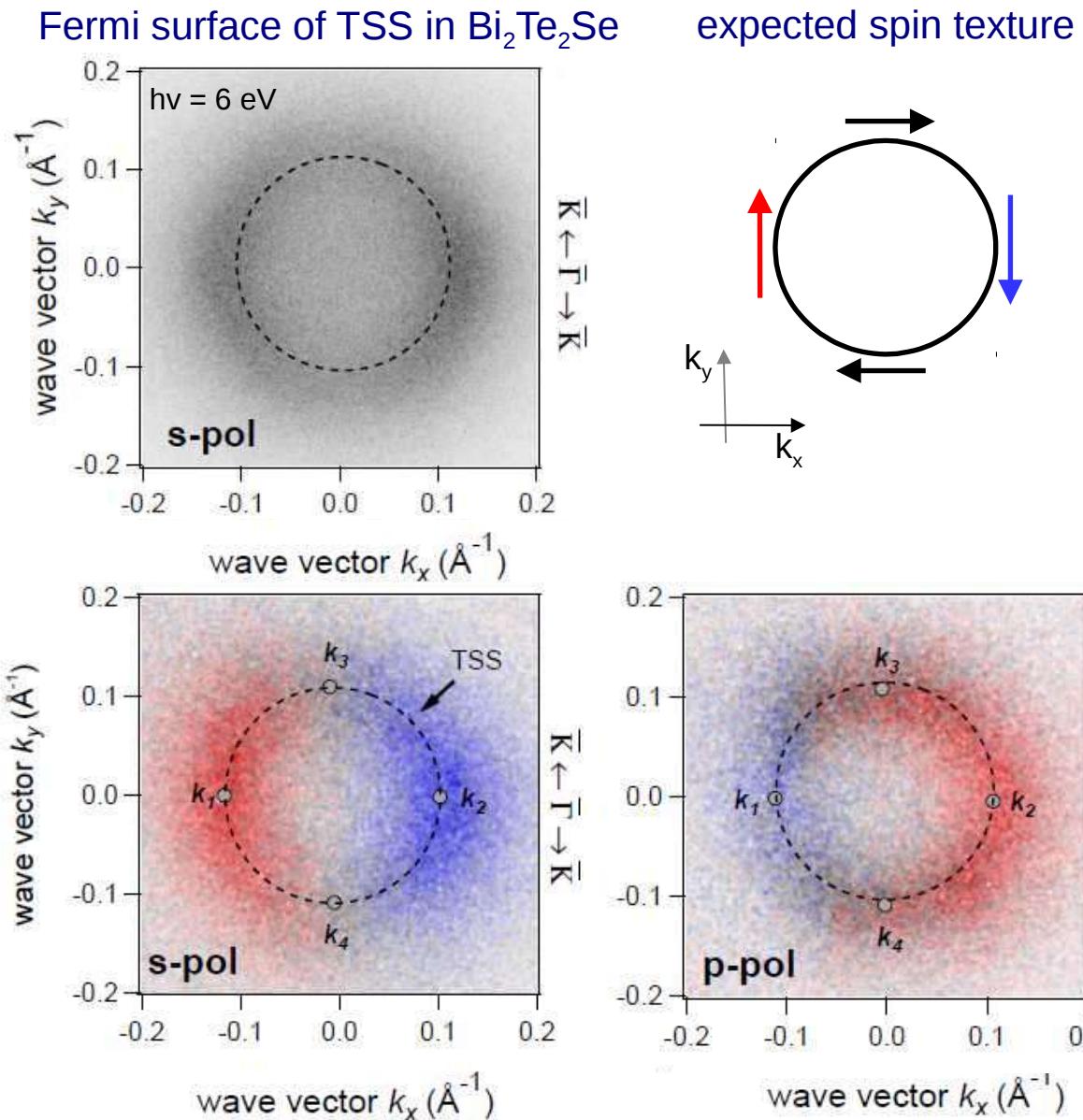
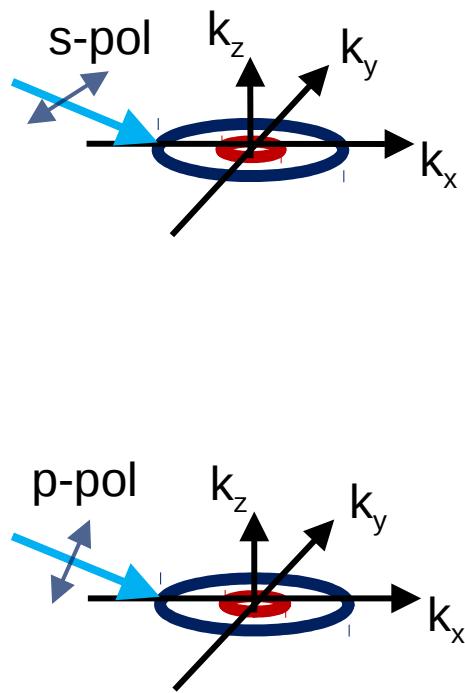
intrinsic spin      ?      photoelectron spin



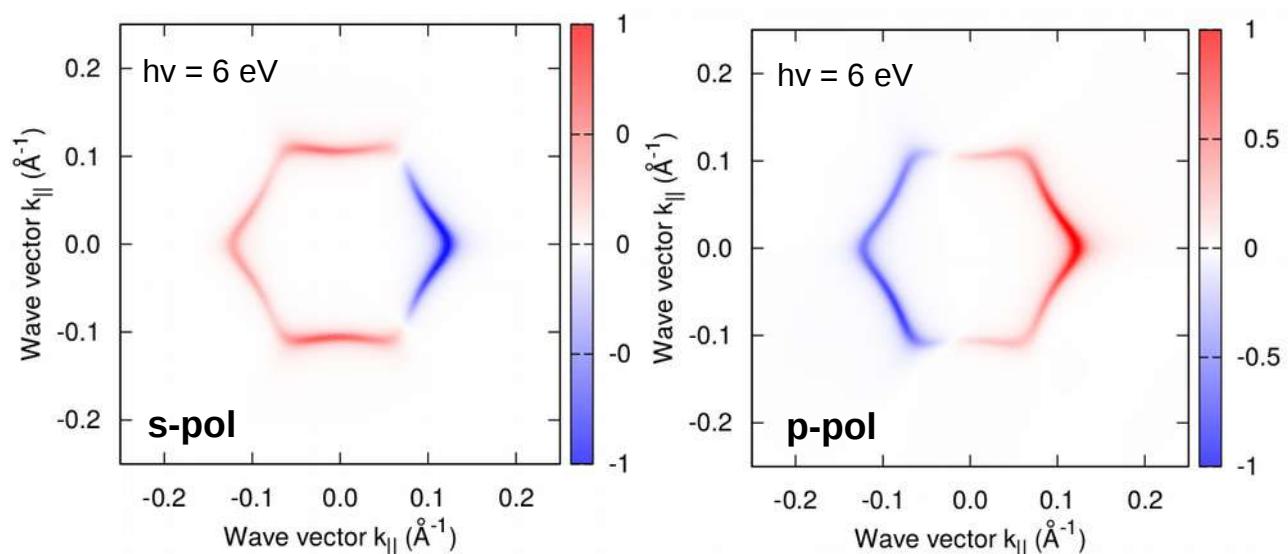
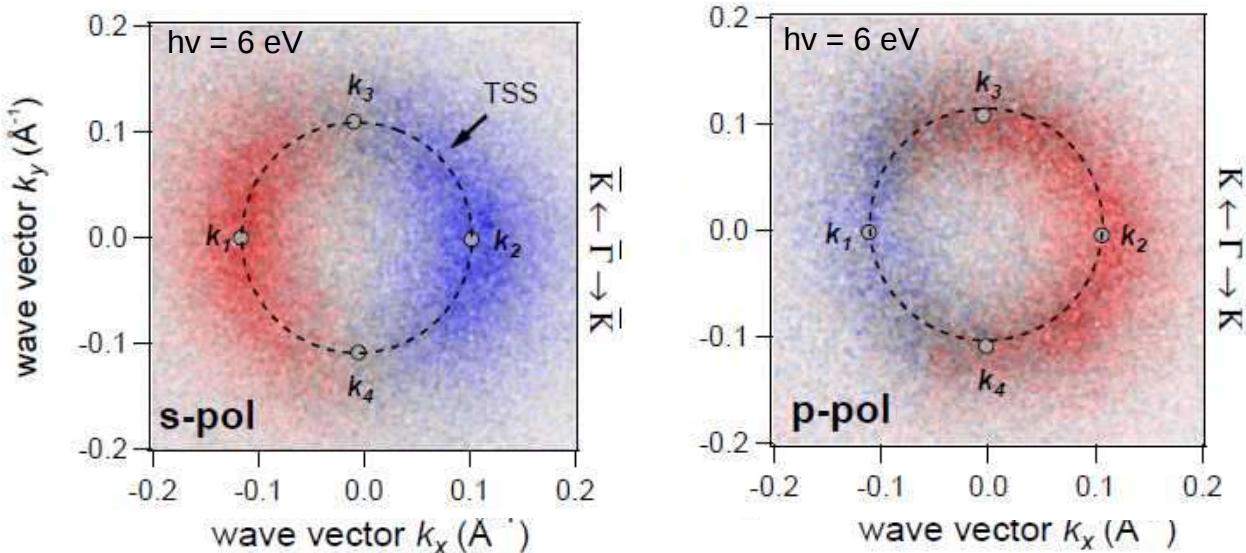
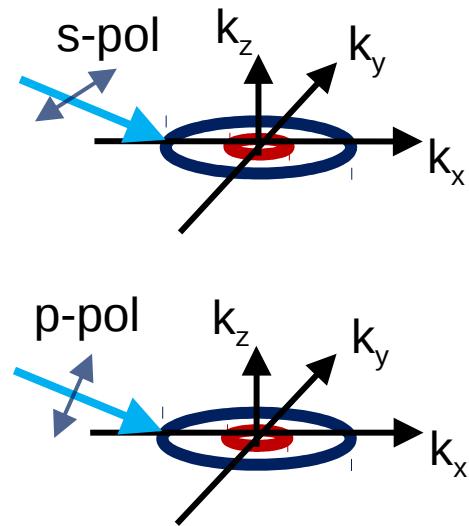
PRL 114 (2015) 066802

- electronic structure probed with simultaneous energy, momentum, and spin sensitivity
- relation between intrinsic and photoelectron spin texture?

# Light-polarization dependence of photoelectron spin



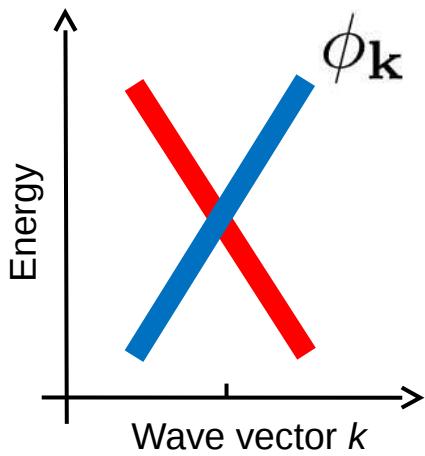
# Experiment vs. one-step photoemission theory



Calculations by Jürgen Braun  
and Jan Minar (LMU, UWB)



# Spin-orbit coupling in Bloch wave functions



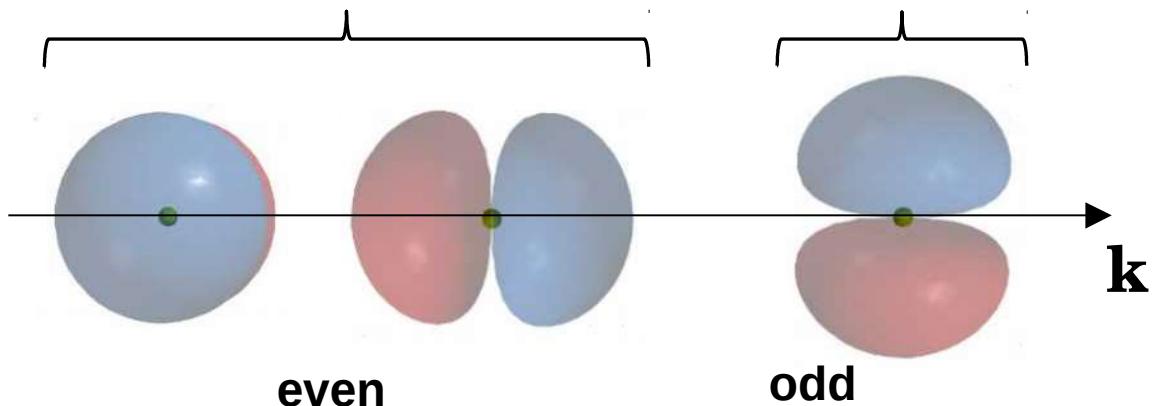
Bloch wave:

$$\phi_{\mathbf{k},\sigma}(\mathbf{r}) = u_{\mathbf{k}}(\mathbf{r})e^{i\mathbf{k}\cdot\mathbf{r}}\chi_{\sigma}$$

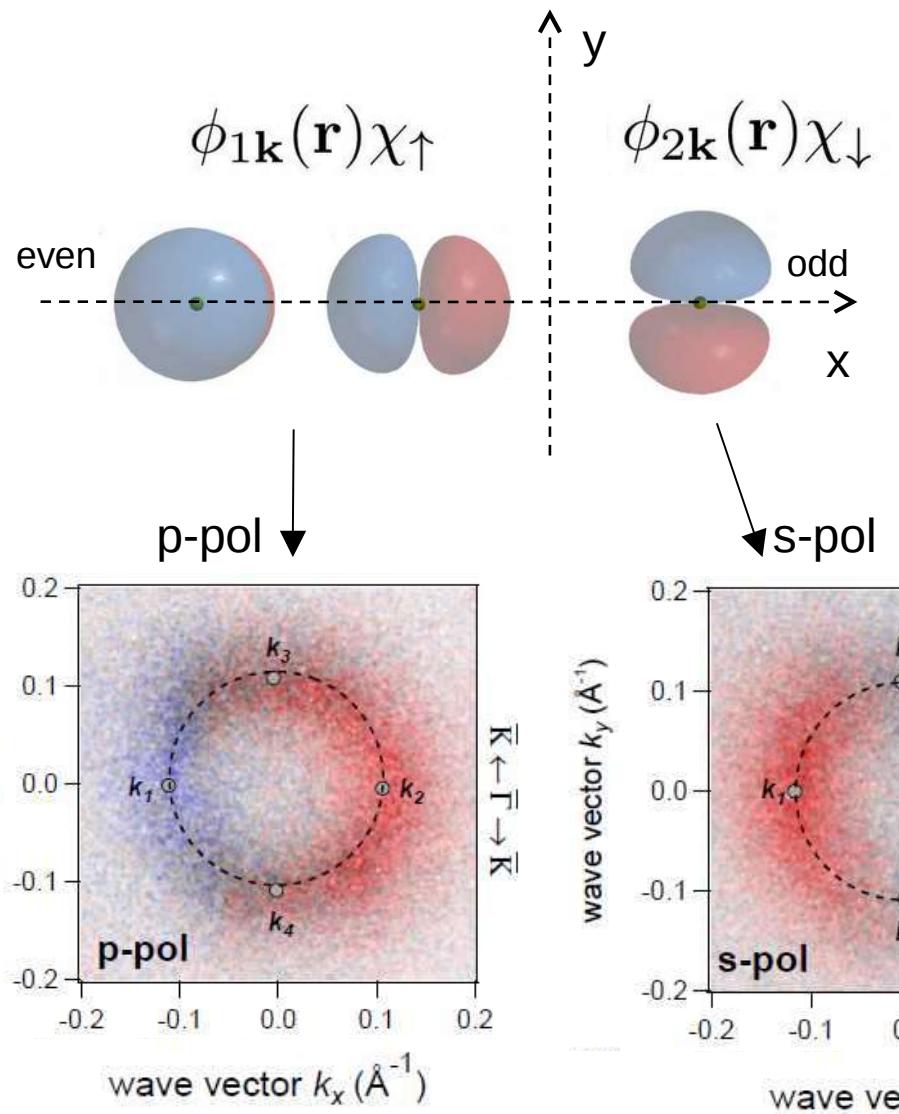
Bloch wave with spin-orbit interaction:

$$\phi_{\mathbf{k}}(\mathbf{r}) = \phi_{1\mathbf{k}}(\mathbf{r})\chi_{\uparrow} + \phi_{2\mathbf{k}}(\mathbf{r})\chi_{\downarrow}$$

when  $\mathbf{k}$  in  
mirror plane:

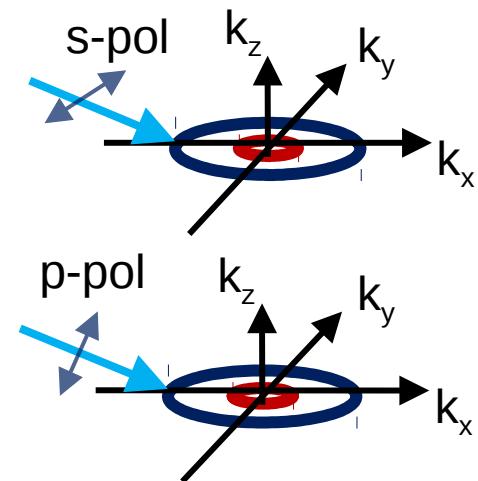


# Dipole selection rules and spin-orbit coupling

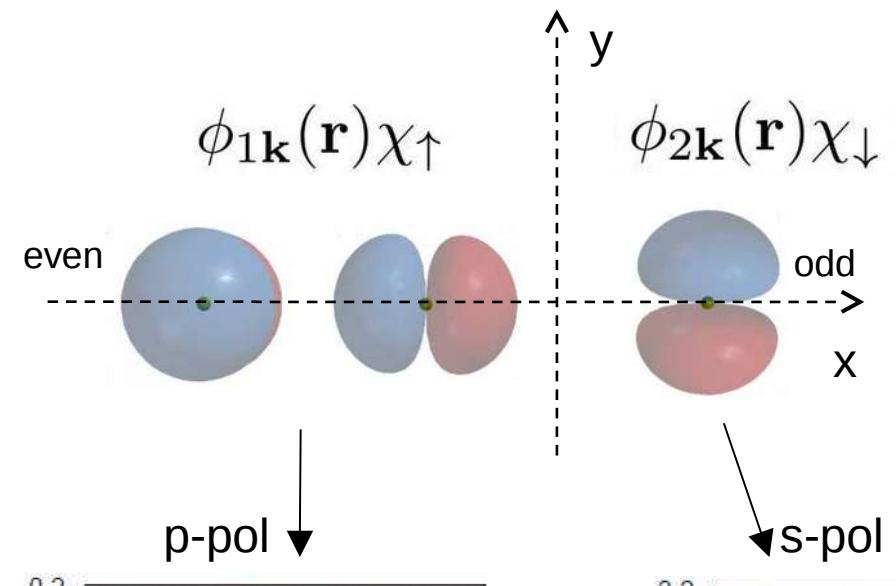


Dipole matrix element:

$$I \propto | \langle \Phi_f | \hat{e} \cdot \mathbf{r} | \phi_{\mathbf{k}} \rangle |^2$$

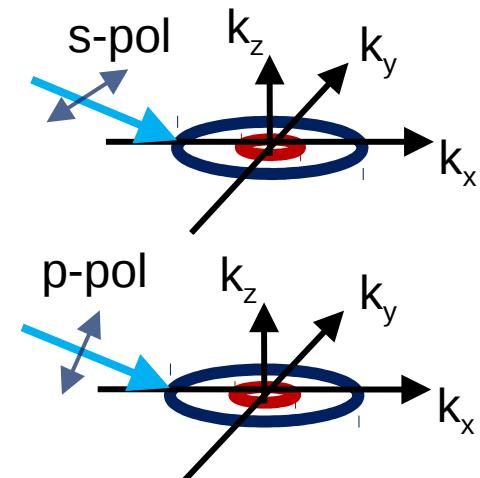


# Dipole selection rules and spin-orbit coupling



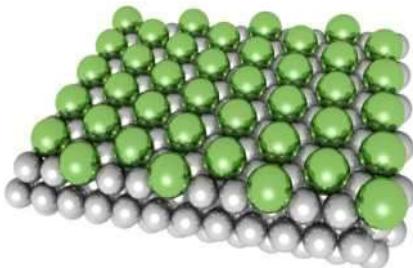
Dipole matrix element:

$$I \propto | \langle \Phi_f | \hat{e} \cdot \mathbf{r} | \phi_{\mathbf{k}} \rangle |^2$$



- photoelectron spin depends on orbital symmetry and experimental geometry
- even the observed spin reversion reflects the **intrinsic spin** structure  
→ talk by Koji Miyamoto for d-states of W(110)

# Model system: Rashba-split surface states on BiAg<sub>2</sub>/Ag(111)



1/3 ML Bi on Ag(111)  
( $\sqrt{3} \times \sqrt{3}$ )R30°-reconstruction

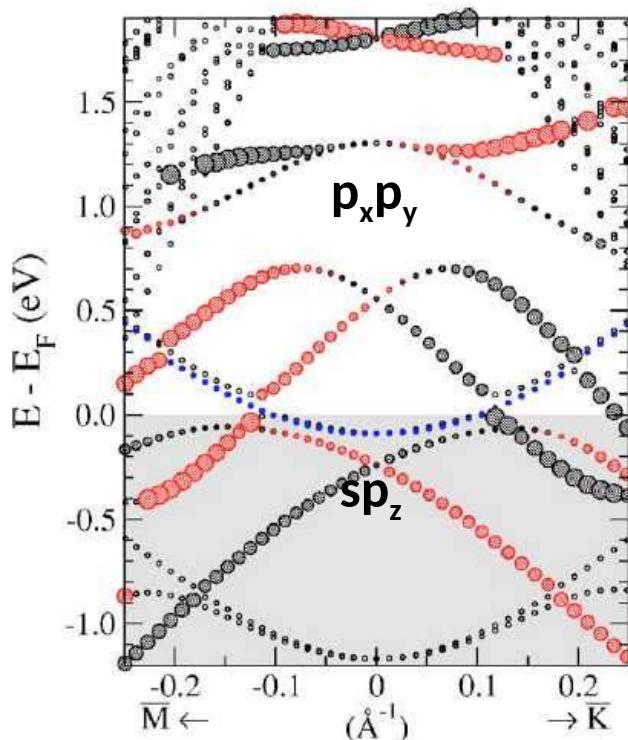


Poster by Ryo Noguchi 1-15

C. Ast *et al.*, PRL **98**,  
186807 (2007)

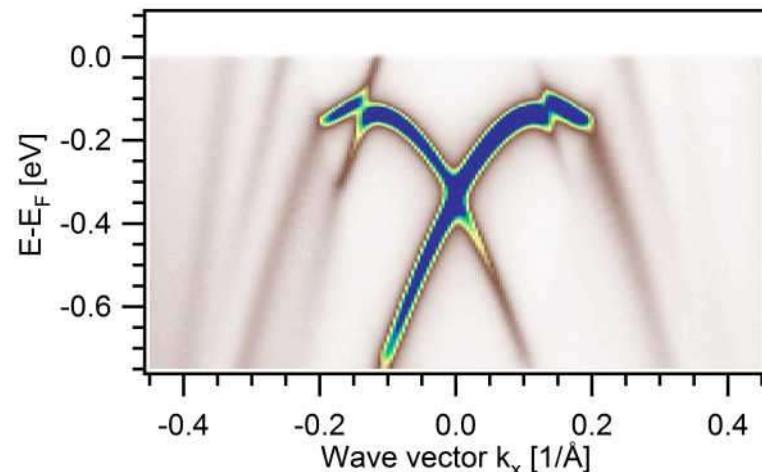
R. Noguchi *et al.*, PRB **95**,  
041111(R) (2017)

## DFT calculation



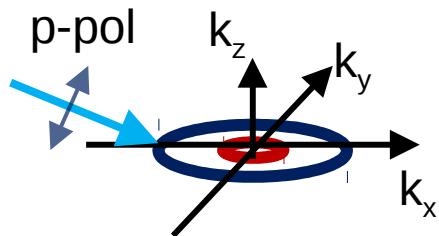
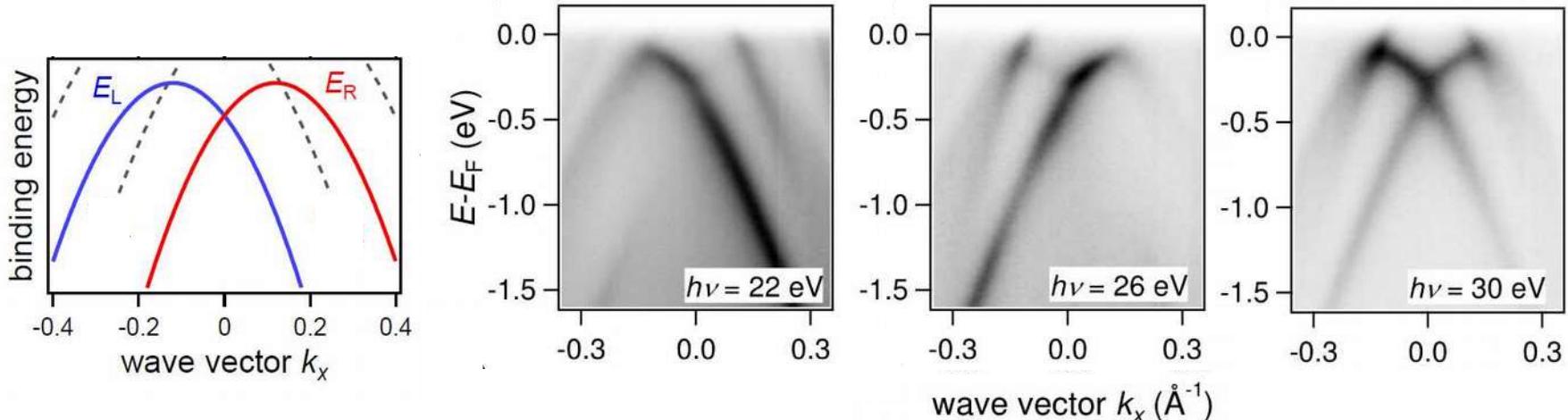
G. Bihlmayer *et al.*, PRB **75**, 195414 (2007)

- nearly perfect 2DEG with
- large spin-orbit splitting  
→ spin-polarized bands
- occupied states of mainly spz character
- unoccupied states of mainly p<sub>x</sub>p<sub>y</sub> character



H. Bentmann *et al.*, EPL **87**, 37003 (2009)

# Linear dichroism for spin-split states in BiAg<sub>2</sub>/Ag(111)



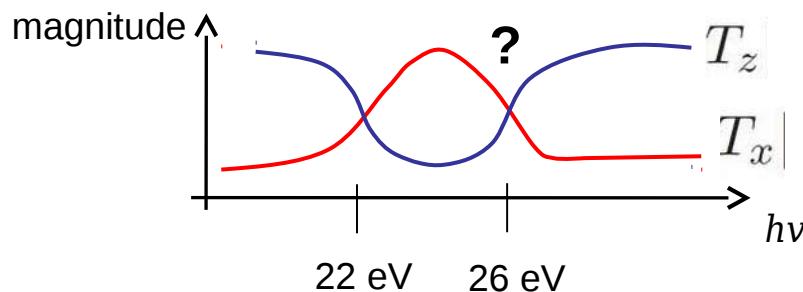
$$I_R \propto | \langle \Phi_f | E_z \hat{z} | \phi_{\mathbf{k}} \rangle + \langle \Phi_f | E_x \hat{x} | \phi_{\mathbf{k}} \rangle |^2$$

$$I_L \propto | \langle \Phi_f | E_z \hat{z} | \phi_{\mathbf{k}} \rangle - \langle \Phi_f | E_x \hat{x} | \phi_{\mathbf{k}} \rangle |^2$$

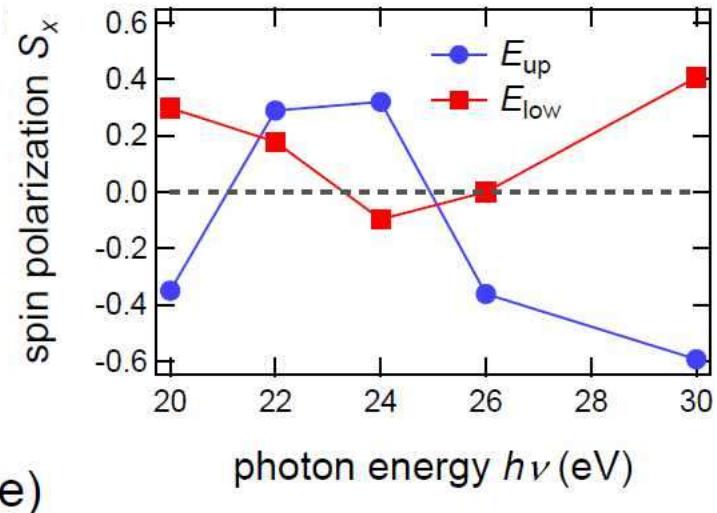
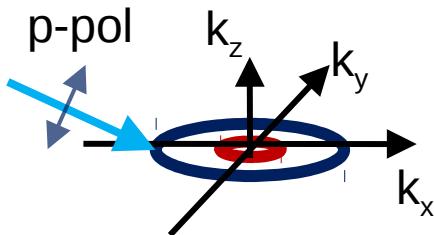
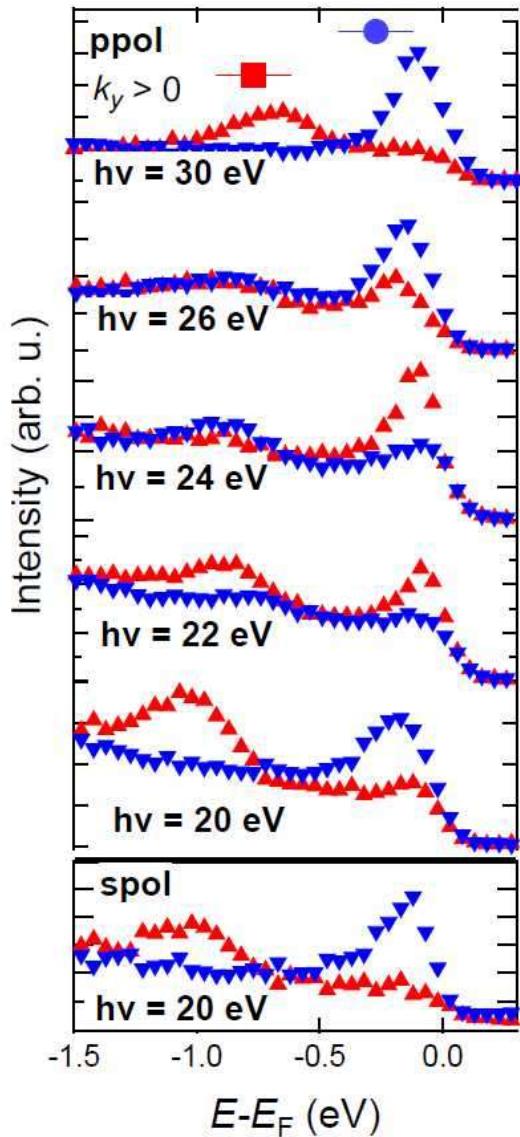
$$\downarrow \quad \quad \quad \downarrow$$

$$T_z \quad \quad \quad T_x$$

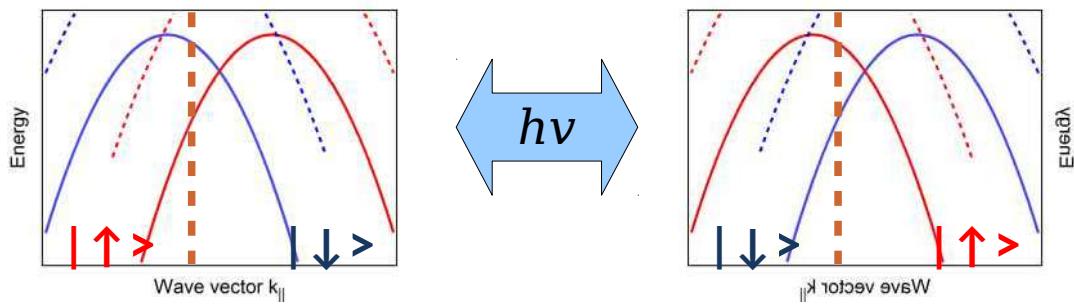
→  $T_x$  and  $T_z$  vary with  $h\nu$ , due to  
change of the final state  $|\Phi_f\rangle$



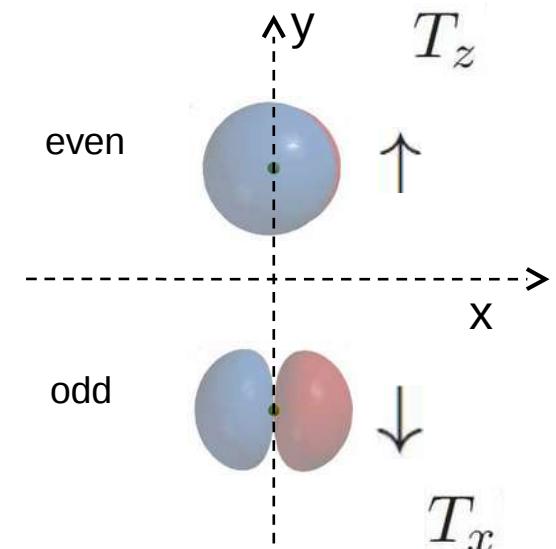
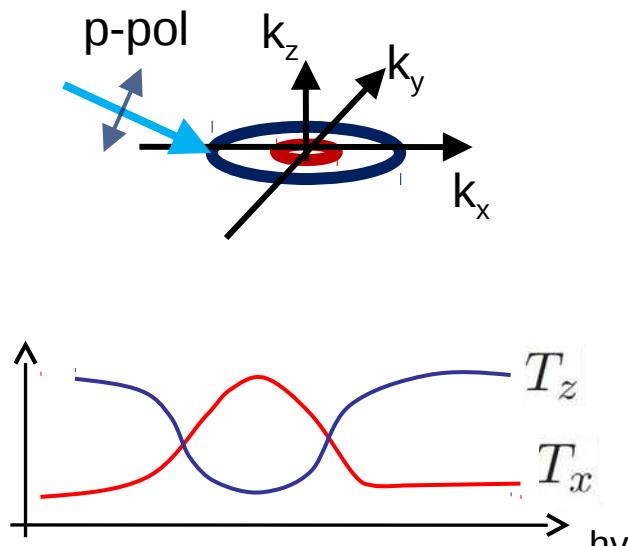
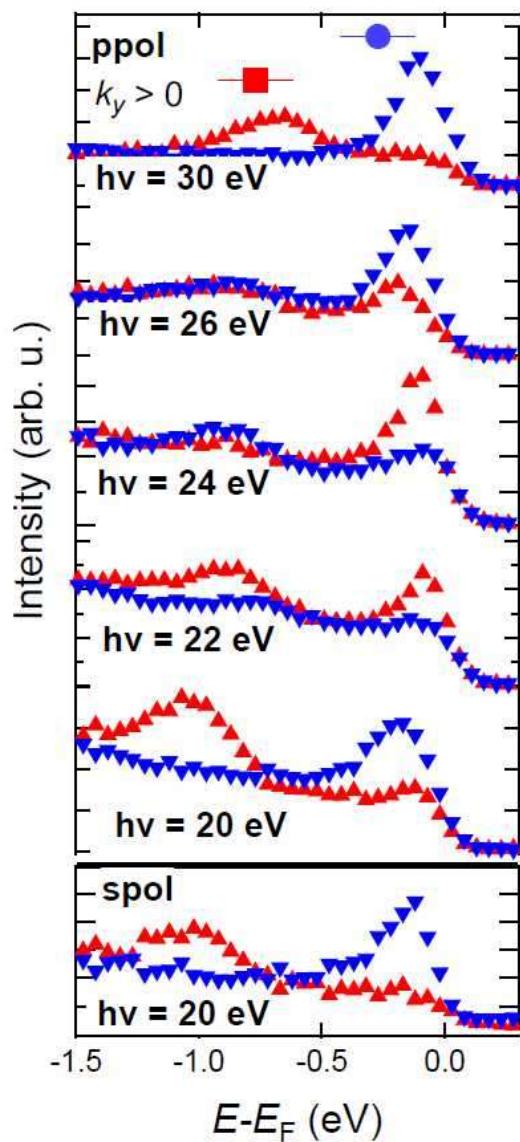
# Photon energy dependence of spin-polarization



→ measured spin-polarization reverses within a few eV of photon energy!

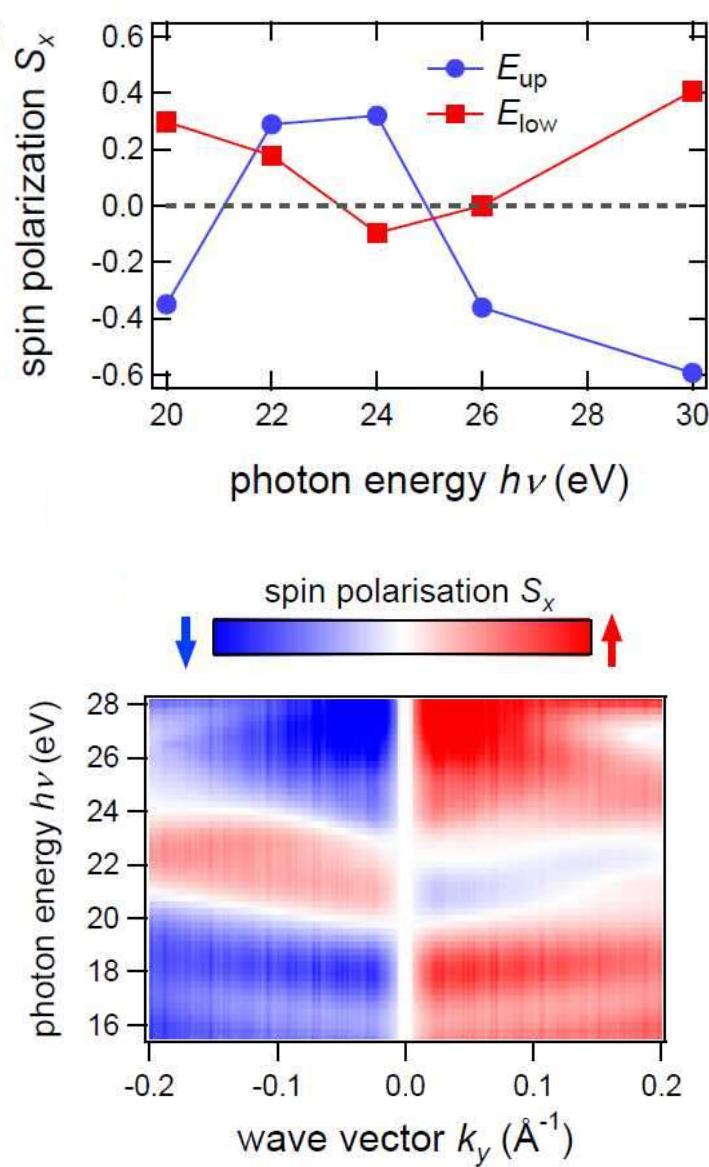
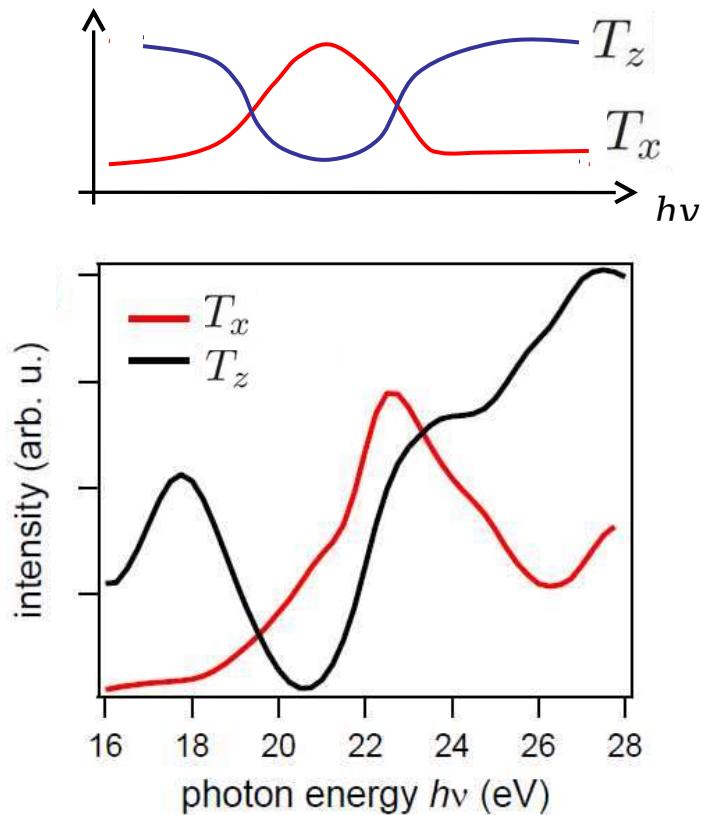


# Photon energy dependence of spin-polarization



- photoelectron spin flips sign with excitation energy
- $T_z$  produces spin- $\uparrow$  and  $T_x$  spin- $\downarrow$  photoelectrons

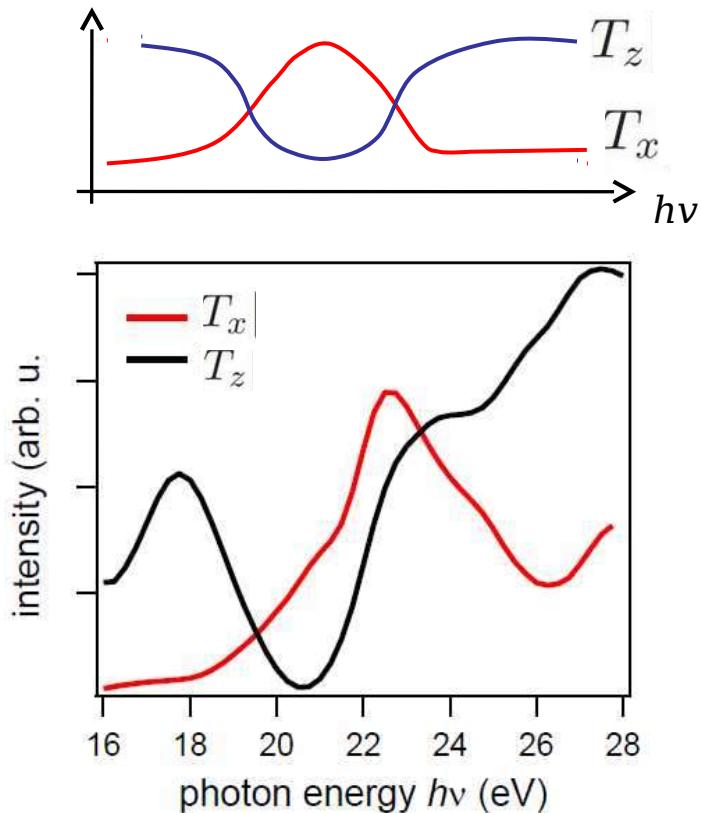
# *Ab initio* one-step photoemission theory



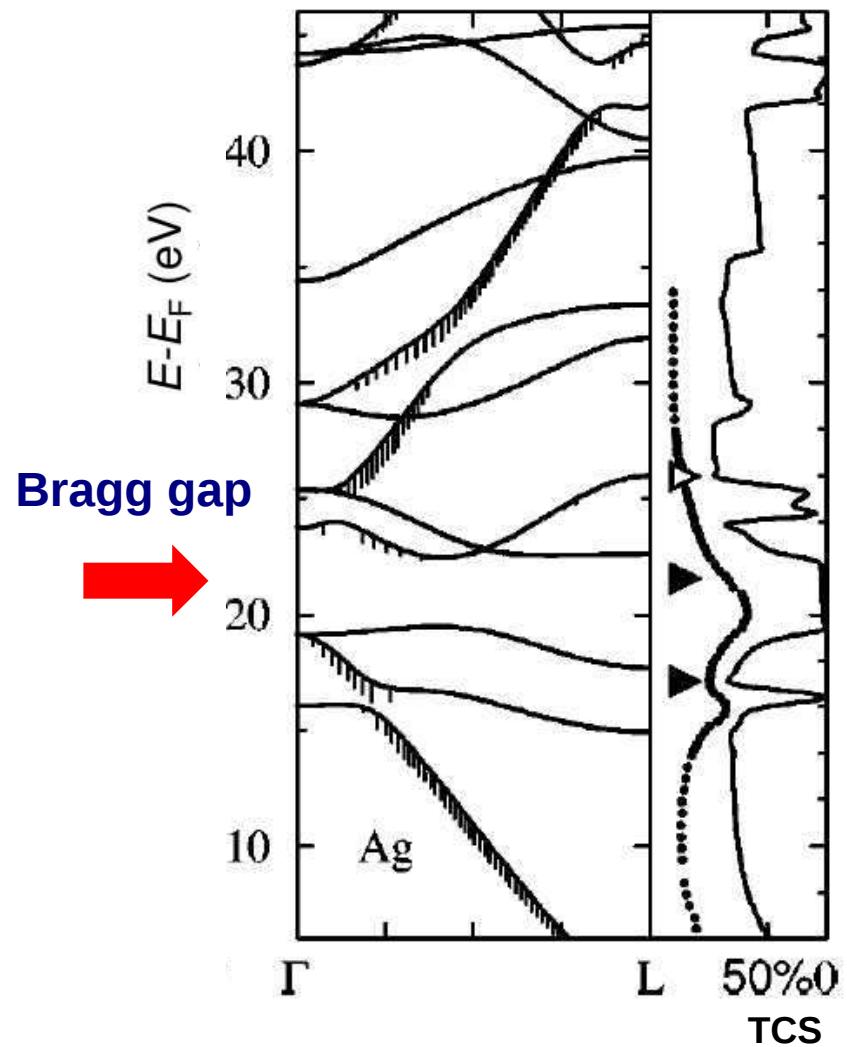
→  $h\nu$ -dependence  
captured by  
*ab initio* theory

E. E. Krasovskii  
DIPC San Sebastian

# *Ab initio* one-step photoemission theory



Ag(111) band structure above  $E_F$



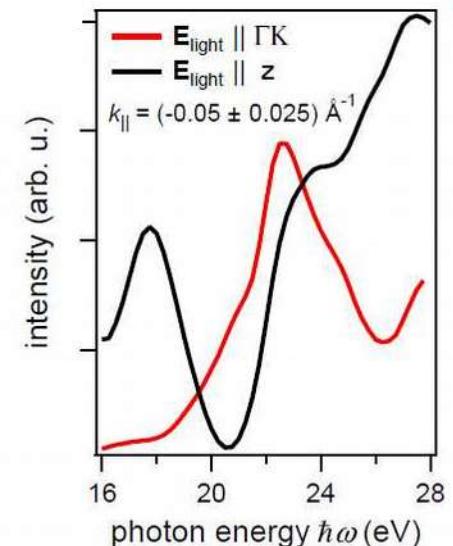
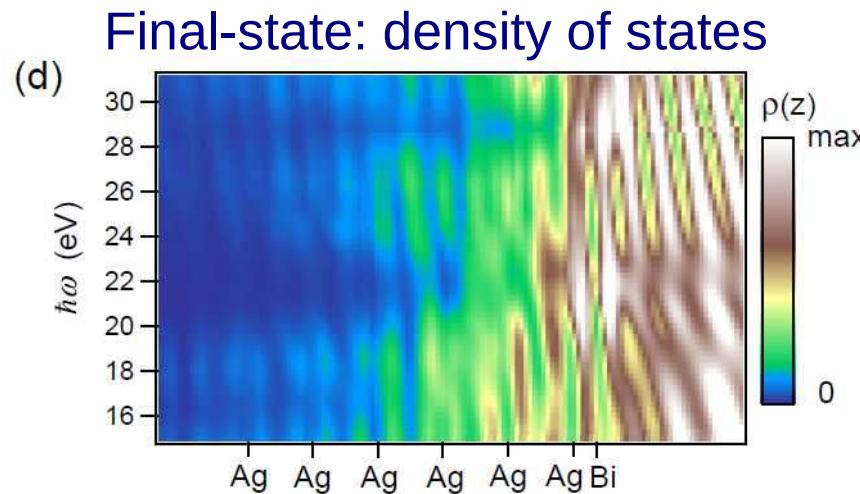
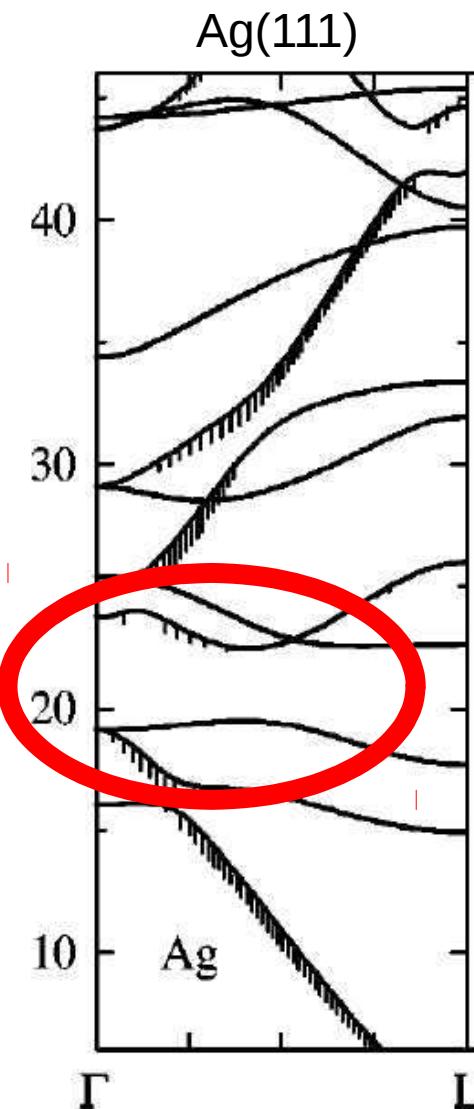
→  $h\nu$ -dependence  
captured by  
*ab initio* theory



E. E. Krasovskii  
DIPC San Sebastian

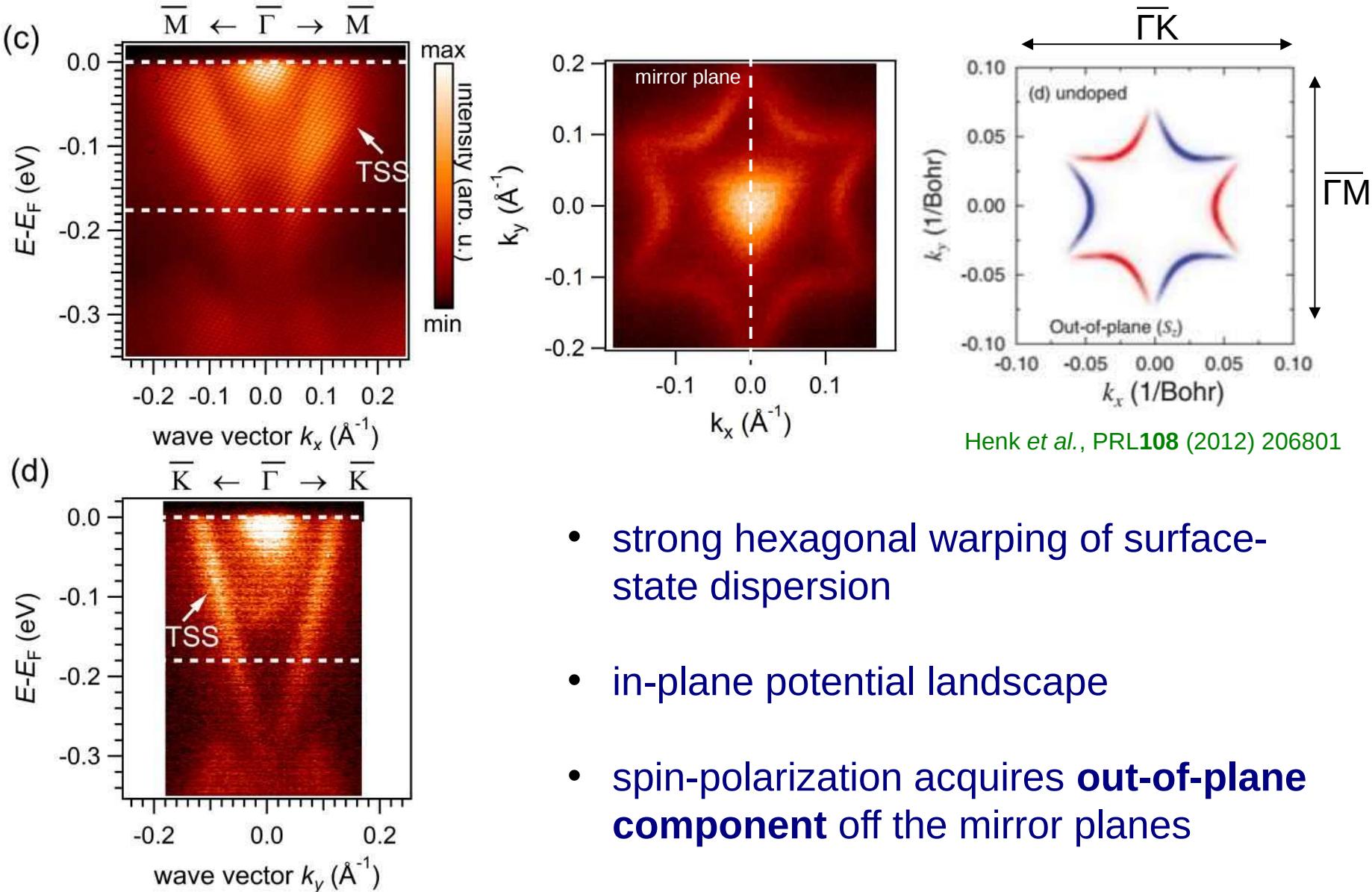
Krasovskii, Schattke, PRB **59** (1999) 15609(R)

# Complex unoccupied band structure: PE Final state

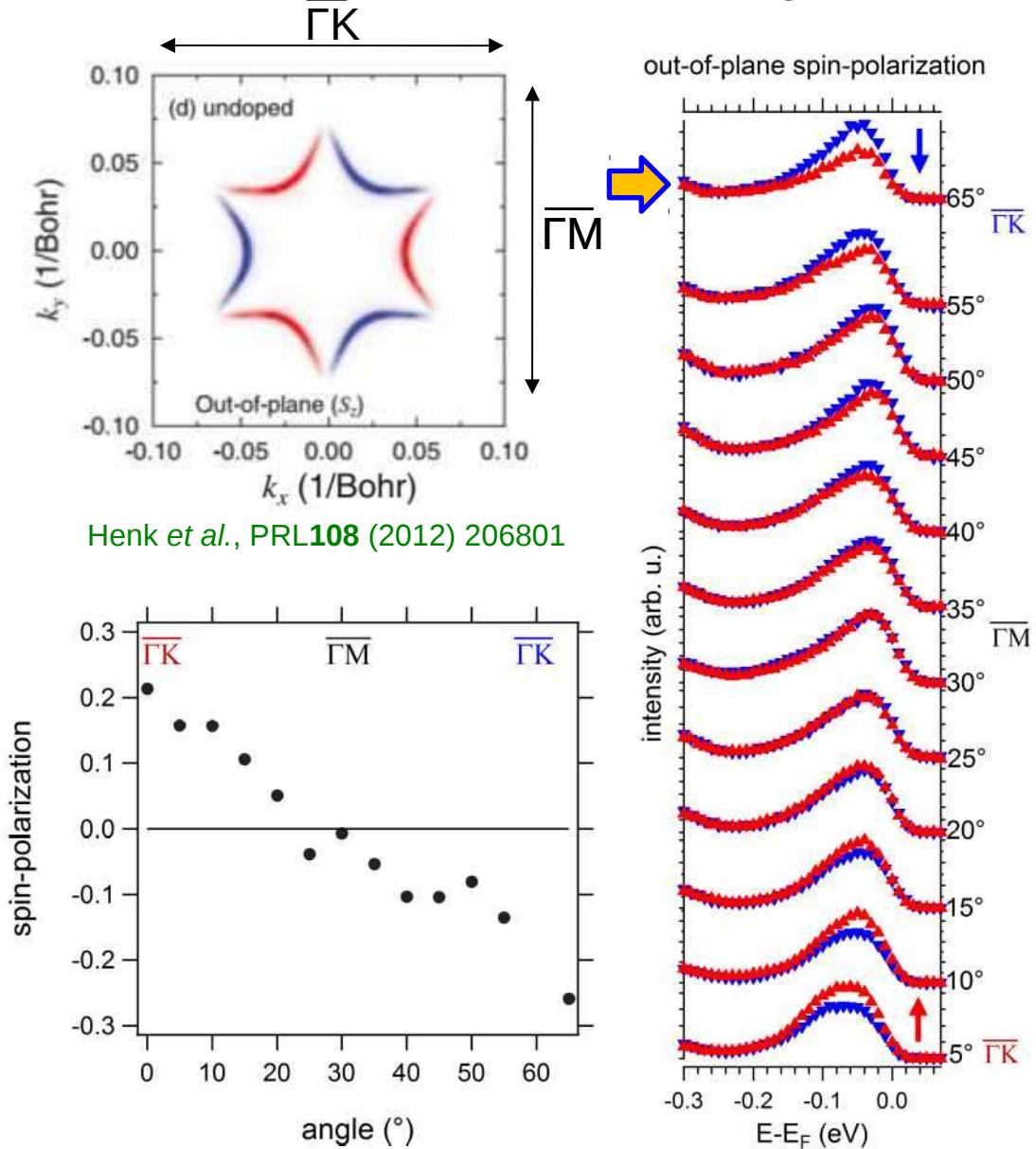
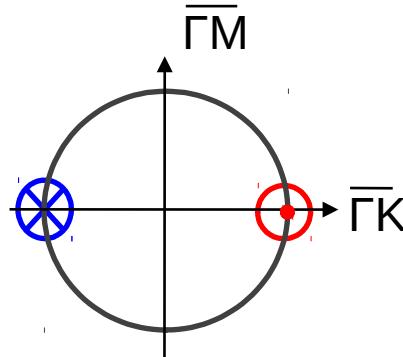
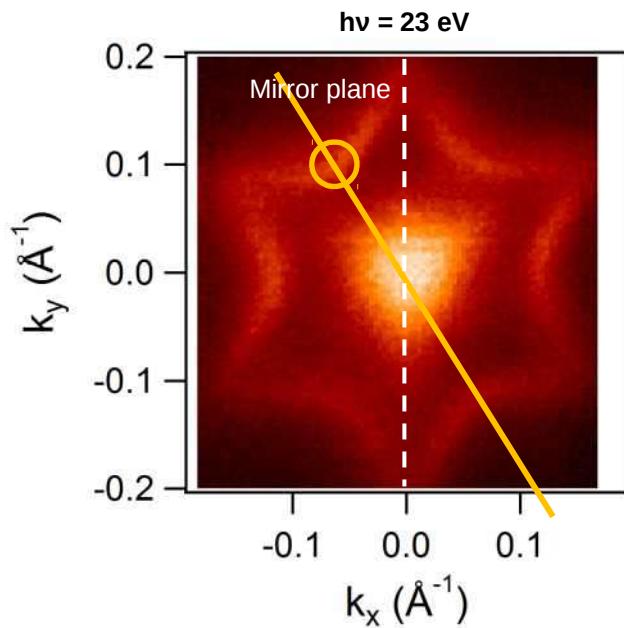


- propagating and evanescent waves
- gap: only evanescent solutions
  - rapid change of final-state wave function

# Topological surface state in $\text{Bi}_2\text{Te}_3$



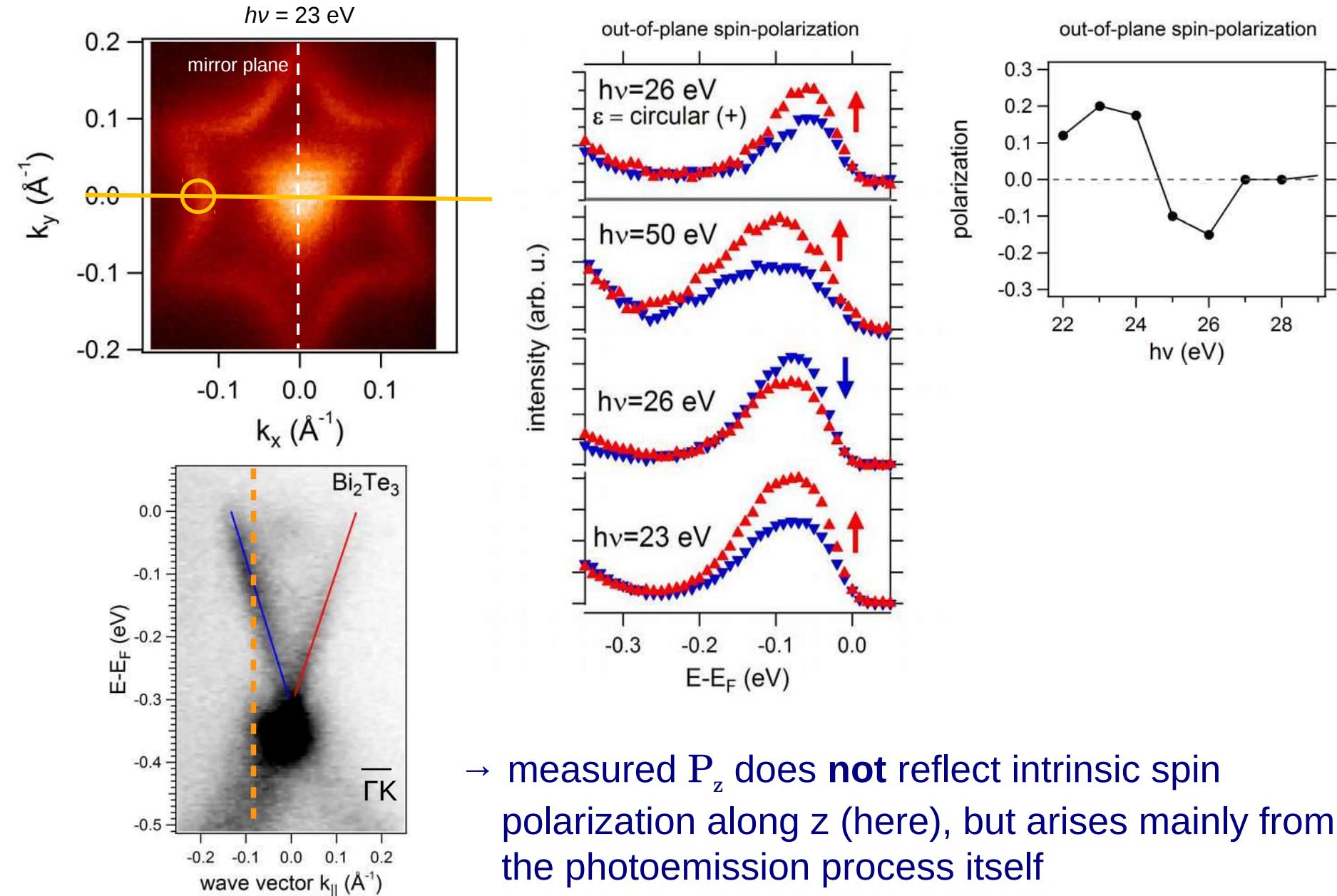
# Out-of-plane spin polarization in $\text{Bi}_2\text{Te}_3$



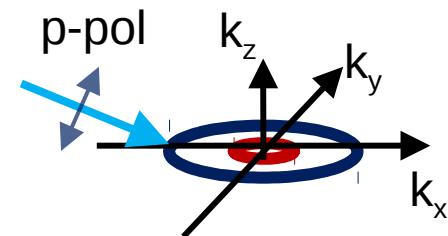
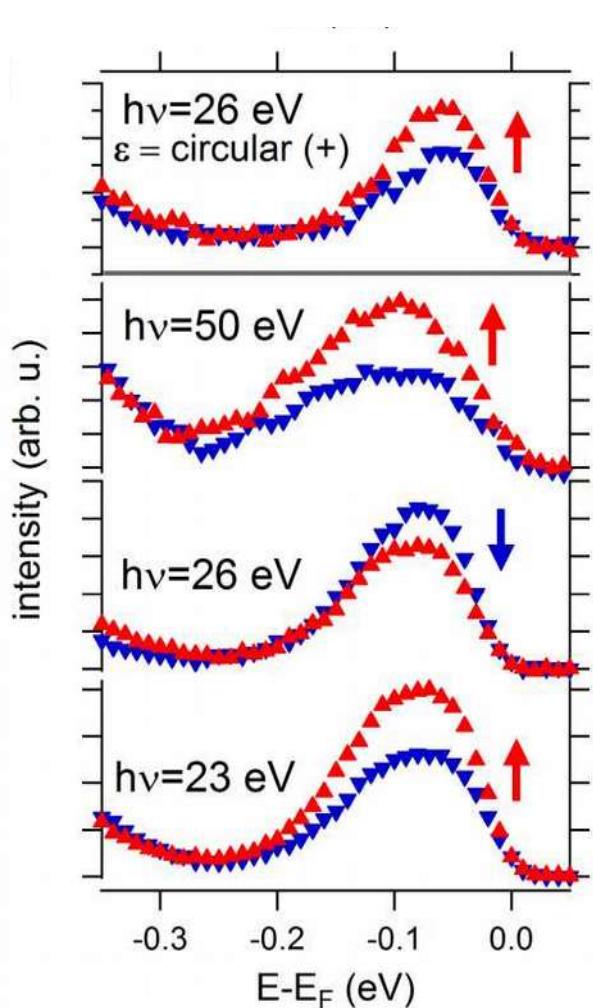
Unpublished: HiSOR (ESPRESSO), Hiroshima

H. Bentman, FR et al., (2017) unpublished

# Effect of photon-energy on out-of-plane spin polarization



# Photon-energy-dependence of out-of-plane spin polarization



Matrix elements:

$$\phi_{\mathbf{k}}(\mathbf{r}) = \phi_{1\mathbf{k}}(\mathbf{r})\chi_{\uparrow} + \phi_{2\mathbf{k}}(\mathbf{r})\chi_{\downarrow}$$

$T_1$        $\Phi_f$        $T_2$

$\mathbf{k}$  along  $\Gamma K$  (no mirror plane)     $T_1, T_2 \neq 0$

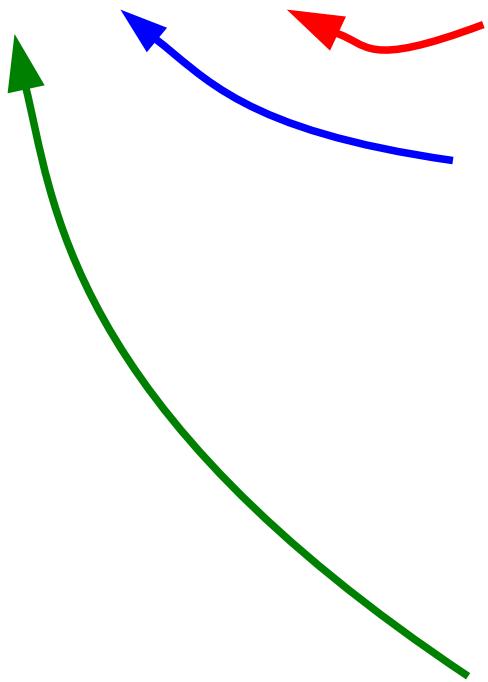
Photoelectron spin:

$$\mathbf{P}_f = \begin{pmatrix} |T_1|^2 - |T_2|^2 \\ |T_1||T_2|\cos(\varphi_1 - \varphi_2) \\ |T_1||T_2|\sin(\varphi_1 - \varphi_2) \end{pmatrix}$$

H. Bentman, FR et al., (2017) unpublished

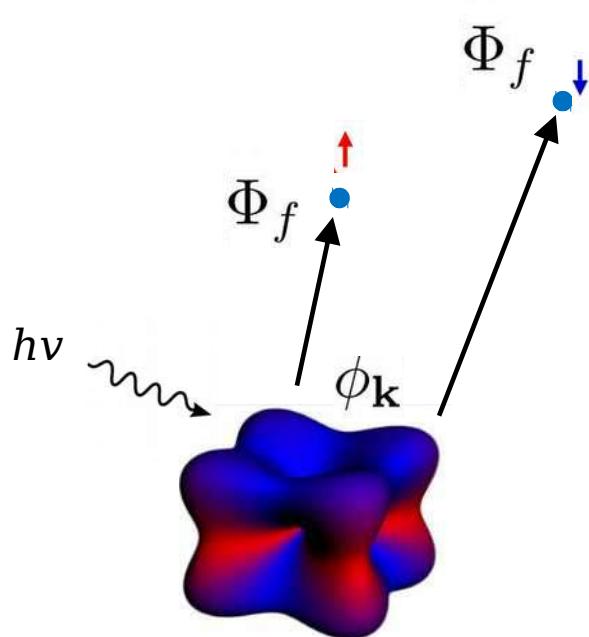
# Summary

$$I \propto | \langle \Phi_f | \hat{e} \cdot \mathbf{r} | \phi_{\mathbf{k}} \rangle |^2$$



- 1) **SOC** couples even and odd WFs differently to the “intrinsic” spin in  $\phi_{\mathbf{k}}$
- 2) **Experimental geometry** (light polarization) “selects” WF and therewith the (photoelectron) spin ( $\rightarrow$  linear dichroism)
- 3) Different contributions do **interfere** in squared matrix element
- 4) If strong **energy dependence** of  $\Phi_f$  (e.g. near gap), interference terms might produce changes of photoelectron spin by changing  $h\nu$

# Conclusions



$$I \propto |\langle \Phi_f | \hat{e} \cdot \mathbf{r} | \phi_{\mathbf{k}} \rangle|^2$$

- Photoelectron spin polarization can strongly depend on experimental parameters
- Spin density of initial state is sampled differently depending on experimental geometry and excitation energy
- In general: comprehensive data and/or additional knowledge from theory or experiment is required
- Outlook: soft X-ray SARPES & experimental info about final states

Fig. from Jackeli et al., Phys. Rev. Lett. 102, 017205 (2009)